# "Meson-scopy": Looking into New Physics via Mesons



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#### **Outline**

#### I. Motivation of Dark Sector Probes

#### II. Meson-scopic Dark-Sector Particle Searches: Intensity Frontier

 Exotic charged meson decays, dark-sector interpretation of the MiniBooNE excess, light dark matter search at SBND

#### III. Meson-scopic Dark-Sector Particle Searches: Energy Frontier

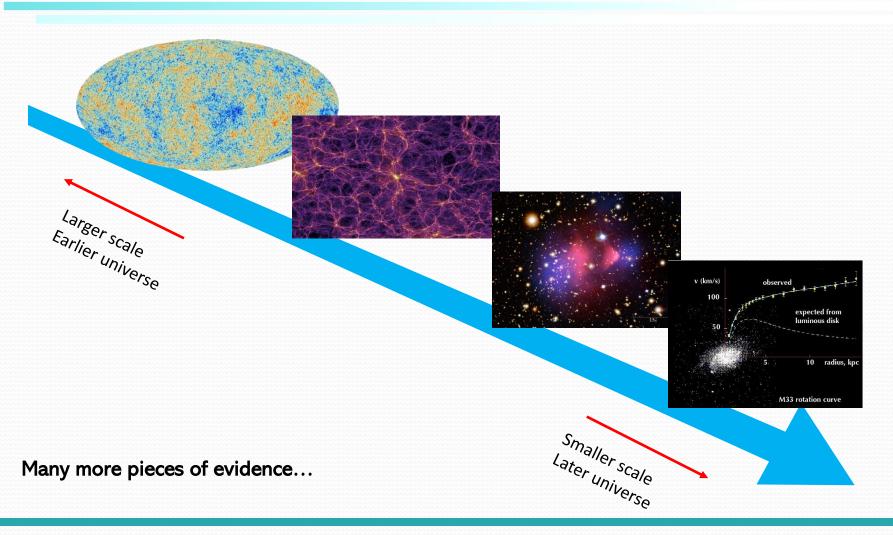
 Search at forward physics facilities (e.g., FASER), search in the LHC detectors (e.g., muon system of CMS)

#### IV. Meson-scopic Dark-Sector Particle Searches: Cosmic Frontier

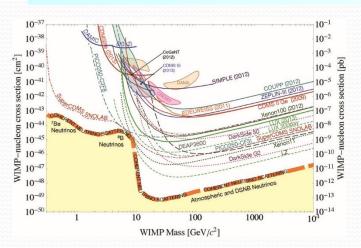
Charged mesons in the cosmic shower

#### V. Conclusions

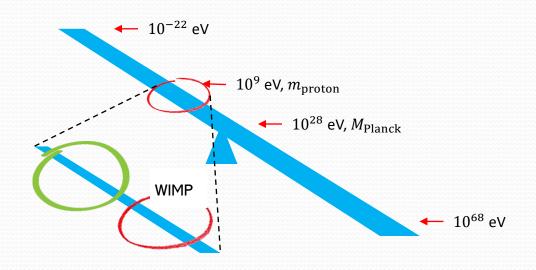
#### **Dark Matter Puzzle**



#### **Light Dark Matter and Light Mediators**



Null observation of (WIMP) dark-matter signal through non-gravitational interactions



#### Light dark matter and light mediators

- Can be a thermal dark matter candidate
- ✓ Less constrained by current searches
- ✓ Often <u>feebly/weakly interacting</u> with SM particles and involving <u>similar mass-range mediators</u>

#### Mediator as a Portal to Dark Sectors

- □ Spin 0
  - Higgs portal  $\mathcal{L}_{\mathrm{Higgs \, portal}}^{d=3,4} = -(\mu S + \lambda S^2) H^{\dagger} H$
- ☐ Spin ½
  - Neutrino portal  $\mathcal{L}_{
    m neutrino\,portal}^{d=4} = -\sum y_
    u^{lpha I}(ar{L}_lpha H) N_I$
- ☐ Spin 1
  - Vector portal  $\mathcal{L}_{ ext{vector portal}}^{d=4} = -rac{\epsilon}{2\cos heta_W} B_{\mu
    u} F'_{\mu
    u}$

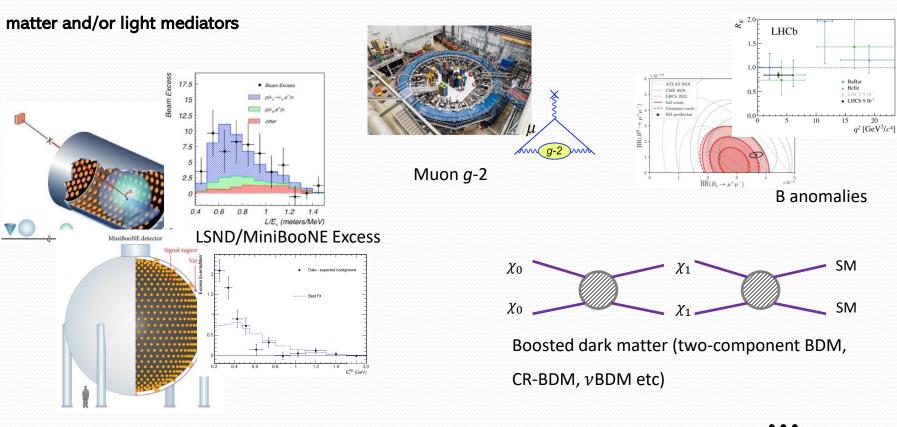


Many models/scenarios have been proposed and related pheno has been investigated for the last decade.

(See also Snowmass2021 Whitepaper "Dark Sector Studies with Neutrino Beams". [arXiv:2207.06898])

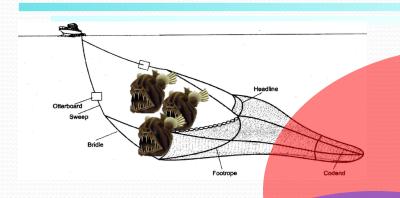
# Various Phenomena and Dark-Sector Interpretations

Various HEP phenomena including excesses, anomalies, non-minimal dark-sector scenarios motivate light dark

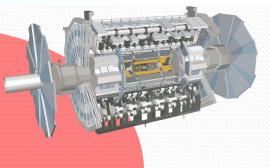


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#### Trawling Dark-Sector Particles from Mesons and Complementarity



**Energy** 





Dark Sector Physics from Mesons

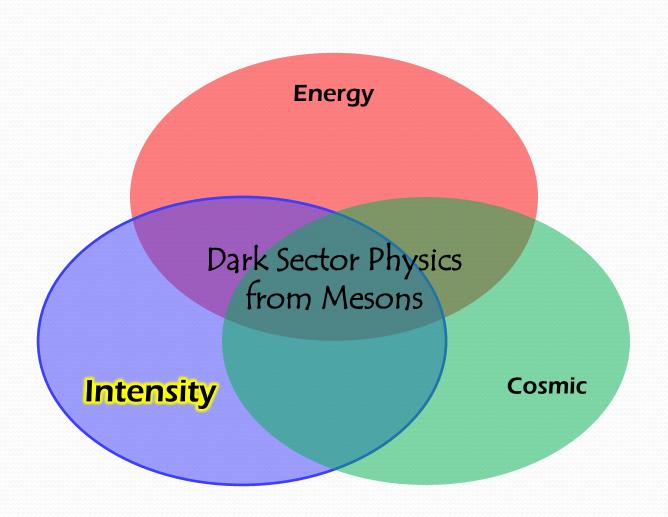


**Intensity** 



Cosmic



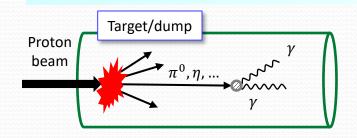


# **NEU**trino facilities are **NEW** physics machines

as well as neutrino factories.

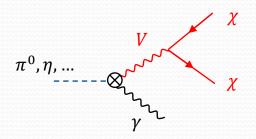
... much more BSM potentials than expected

# Conventional Meson-Induced New Physics Scenarios



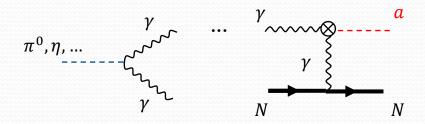
Promptly decaying  $\pi^0$  ,  $\eta$ 

Example 1: vector-portal dark matter  $(\chi)$ 



+ beam bremsstrahlung, ...

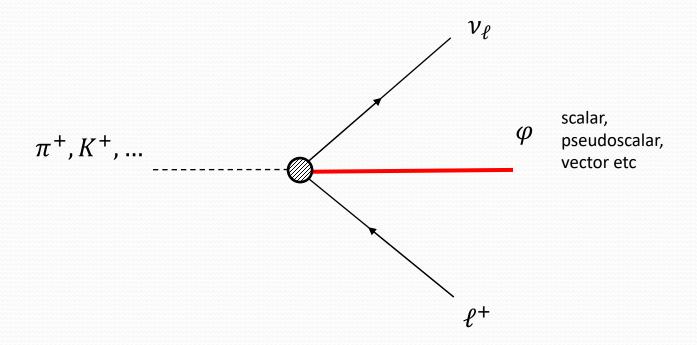
Example 2: axion-like particles (a) interacting with photons



+ E&M cascade shower photons, ...

Neutral mesons have played a crucial role in production of new physics particles.

# **Contributions from Charged Mesons?**



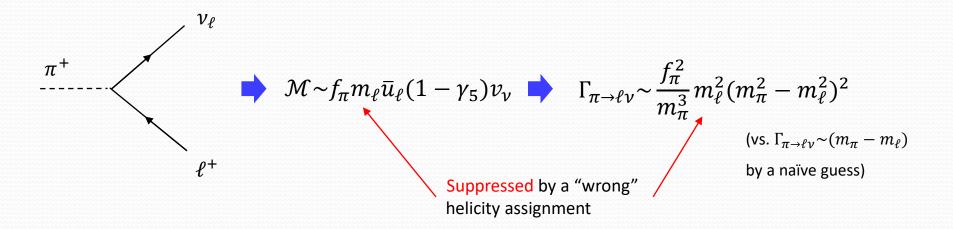
Charged mesons can be focused by a magnetic horn system,

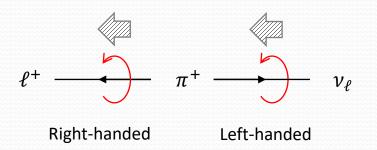


but could this decay be sizable?



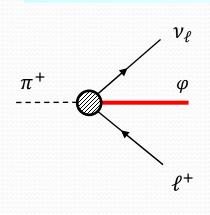
# Two-Body Decay of a Charged Meson

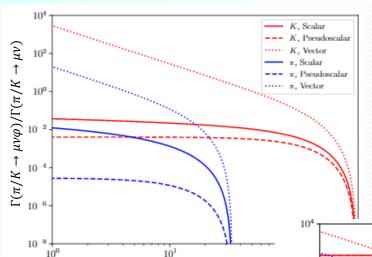




Angular momentum conservation highly suppresses the decay of scalar mesons in this way.

# Three-Body Decay of a Charged Meson

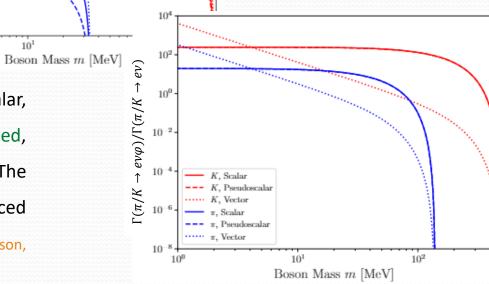




•  $\varphi$  is assumed to couple to the charged lepton only and the associated couplings are set to be unity for comparison.

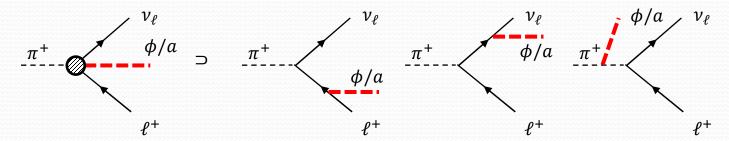
By adding the third particle  $\varphi$  (scalar, pseudoscalar, vector, etc), the helicity suppression can be evaded, i.e., 3-body decays can be **hugely enhanced**. The decay to a massive vector is even more enhanced due to the longitudinal polarization. [e.g., Carlson,

Rislow, arXiv:1206.3587]

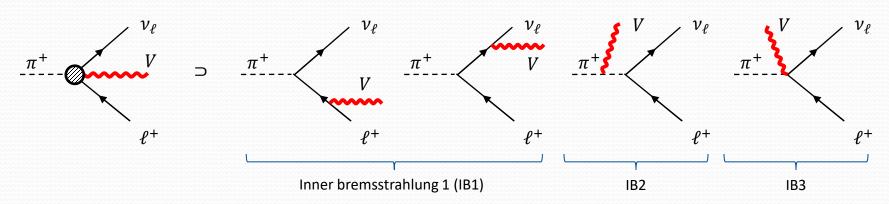


# Various Dynamics in the Three-Body Decay

#### (Pseudo)scalar case



#### Vector case



- For models having couplings to the quark contents inside the meson, (QCD-origin) structure-dependent (SD) terms may arise.
- Typically, IB3 contributions ≫ IB2 ≈ IB1

# Pion Dark Radiative Decay

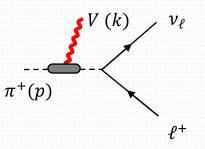
$$\begin{split} \mathcal{M} &= i g_\pi \varepsilon^\mu \left[ \bar{u}_\ell \gamma_\rho (1-\gamma_5) v_\nu \right] \left[ \int \! d^4 x e^{iqx} \langle 0 \, | \, T \{ j_\mu(x) \bar{d} \, \Gamma^\rho u(0) \} \, | \, \pi^+ \rangle \right] \\ j_\mu &= \bar{q} \gamma_\mu q \end{split}$$

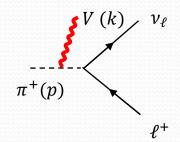
The hadronic part of the matrix element can be decomposed into a generic gauge-invariant form:

$$T_{\mu\rho} = \alpha g_{\mu\rho}(p-k) \cdot k + \beta (p-k)_{\mu} k_{\rho} + i [g_{\mu\rho}(p-k) \cdot k - (p-k)_{\mu} k_{\rho}] F_{A} + \epsilon_{\rho\mu\lambda\sigma} p^{\lambda} k^{\sigma} F_{V} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\rho}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\mu}}{2p \cdot k - m_{X}^{2}} f_{\pi} + \frac{i (2p-k)_{\mu} (p-k)_{\mu}}{2p \cdot$$

Contact interaction

Structure-dependent contributions





IB2

If we wish to satisfy the Ward Identity for our dark current, we require  $\alpha + \beta = \frac{if_{\pi}}{(p-k) \cdot k}$ .

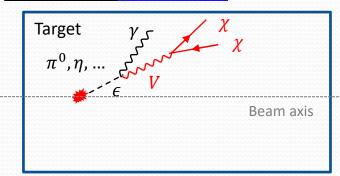
We can choose a family of (unknown) solutions for  $\alpha$  and  $\beta$ .

See also more detailed discussions in [Khodjamirian, Weiler, arXiv:hep-ph/0111249].

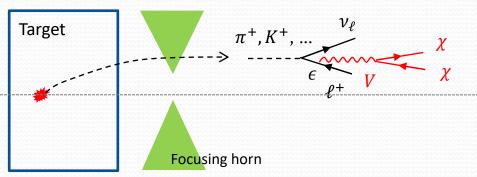
#### Charged Meson vs. Neutral Meson

In the example of vector-portal dark matter,

#### Production via neutral meson



#### Production via charged meson



$$N_{\pi^0-\text{ind}}^{\chi} = N_{\pi^0} \cdot \text{BR}(\pi^0 \to V \to \chi) \cdot f_{\pi^0-\text{ind}}^{\chi}$$

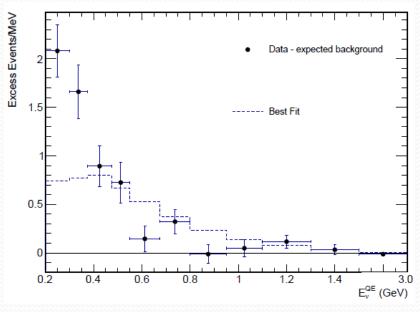
$$N_{\pi^{0}-\mathrm{ind}}^{\chi} = N_{\pi^{0}} \cdot \mathrm{BR}(\pi^{0} \to V \to \chi) \cdot f_{\pi^{0}-\mathrm{ind}}^{\chi} \qquad \langle \langle N_{\pi^{\pm}-\mathrm{ind}}^{\chi} = N_{\pi^{\pm}-\mathrm{ind}}^{\chi} = N_{\pi^{\pm}-\mathrm{ind}}^{\chi} = N_{\pi^{\pm}-\mathrm{ind}}^{\chi}$$

- Comparable production rate:  $\pi^0$ :  $\pi^+$ :  $\pi^-$ :  $\eta$ :  $K^+$ :  $K^- \approx 1$ : 1: 1: 0.1: 0.1: 0.1  $\Rightarrow N_{\pi^0} \approx N_{\pi^\pm}$
- No BR enhancement vs. Large BR enhancement  $\Rightarrow$  BR $(\pi^0 \to V \to \chi) \ll$  BR $(\pi^{\pm} \to V \to \chi)$
- Unfocused  $\pi^0$ ,  $\eta$  vs. Focused  $\pi^{\pm}$ ,  $K^{\pm}$
- Wider spreading  $\pi^0$ ,  $\eta$ -induced flux vs. Forward-directed  $\pi^\pm$ ,  $K^\pm$ -induced flux  $\Rightarrow f_{\pi^0-\text{ind}}^\chi < f_{\pi^\pm-\text{ind}}^\chi$

Production via charged meson can be more efficient than production via neutral meson!

# Example Physics Case 1: Dark-Sector Solutions to the MiniBooNE Low-Energy Excess

#### MiniBooNE (MB) Low Energy Excess

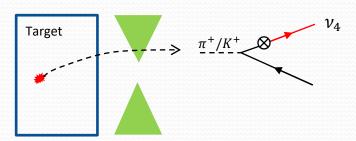


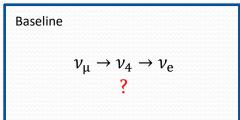
[MiniBooNE Collaboration, arXiv:2006.16883]

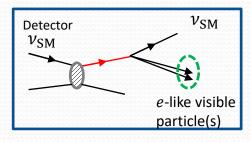
- $\Box$  An observational motivation of new physics (4.8 $\sigma$ )
- Numerous explanations, all involving

  neutrino-sector new physics [Karagiorgi, Djurcic,
  Conrad, Shaevitz, Sorel (2009); Collin, Arguelles, Conrad,
  Shaevitz (2016); Giunti, Lavender (2011); Gariazzo, Giunti,
  Lavender, Li (2017); Kopp, Maltoni, Schwetz (2011); Doring,
  Pas, Sicking, Weiler (2018); Dutta, Ghosh, Li (2020), and
  many more]

# **Conventional Explanations**







New particles in the beam

Sterile neutrino decay

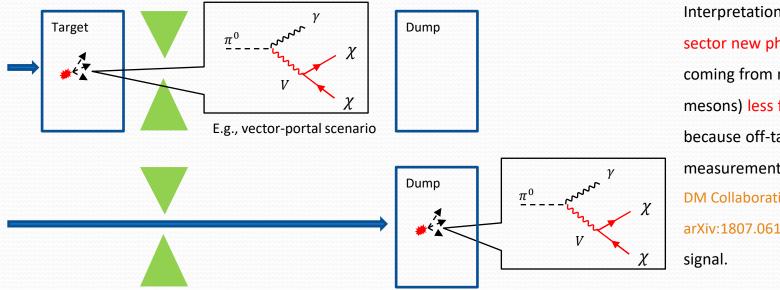
Non-standard flavor transitions

- Oscillations (sterile neutrinos, Lorentz violation, matter effects, large extra dimension etc)
- Lepton number violating decay

New particles produced by neutrino scattering

- Neutrino upscattering involving collimated an  $e^-e^+$  pair
- Neutrino upscattering involving a photon

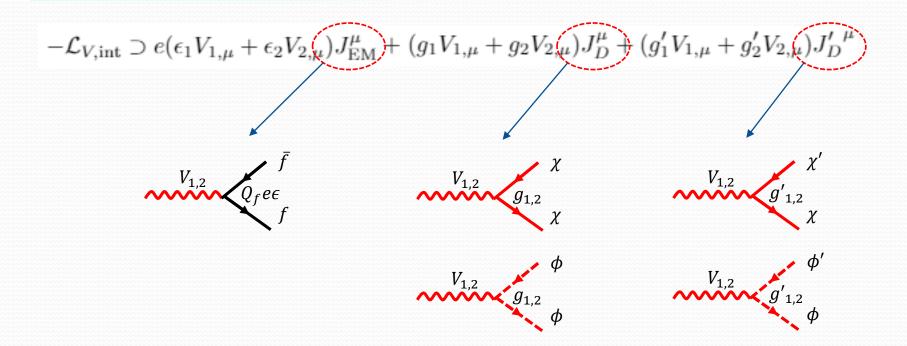
# Dark-Sector Solutions? Challenges?



Interpretations with darksector new physics (mostly coming from neutral mesons) less favored because off-target mode measurements [MiniBooNE DM Collaboration, arXiv:1807.06137] show null

We propose an idea of plausible dark-sector scenarios for the MB excess.

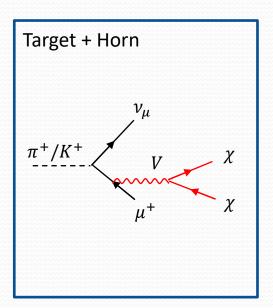
#### An Example Model: Vector-Portal Dark Matter



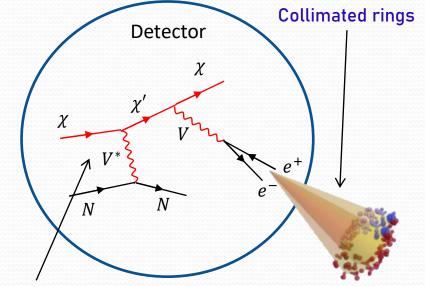
- ☐ Dark photon mediators for purposes of illustration.
- Two different mediators for purposes of generality.
- $\square$  Heavier dark-sector state  $\chi'$ ,  $\phi'$  in addition to dark matter state  $\chi$ ,  $\phi$ : up-scattering  $\chi \to \chi'$ , decay  $\chi' \to \chi$  allowed.

# A Dark-Matter Interpretation for the MB Excess: Main Idea

$$-\mathcal{L}_{V,\text{int}} \supset e(\epsilon_1 V_{1,\mu} + \epsilon_2 V_{2,\mu}) J_{\text{EM}}^{\mu} + (g_1 V_{1,\mu} + g_2 V_{2,\mu}) J_D^{\mu} + (g_1' V_{1,\mu} + g_2' V_{2,\mu}) J_D^{\mu}$$



BR(
$$V \to 2\chi$$
):BR( $V \to 2e$ )  
= 50%:50%



 $BR(V \to 2\chi):BR(V \to 2e) = 50\%:50\%$ 

for illustration

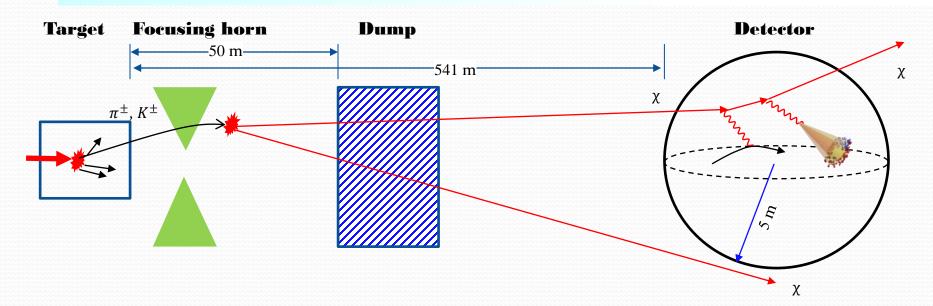
Scattering may happen through an exchange of a different mediator.

[The different mediator could be massless, e.g., dipole operator,  $\overline{\chi'}\sigma^{\mu\nu}\chi F_{\mu\nu}$ ]

Cf. Subdominant  $\pi^0$ -origin contributions are included.

[Dutta, **DK**, Thompson, Thornton, Van de Water, arXiv:2110.11944]

#### Simulation

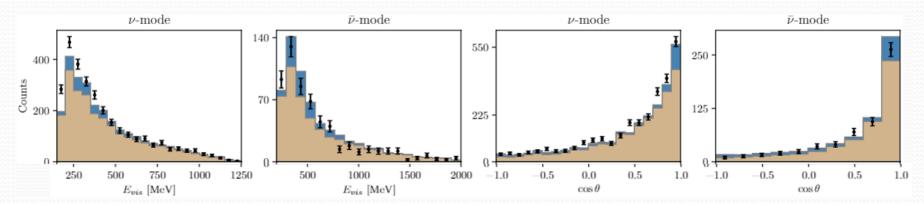


- Sanford-Wang, BMPT-like models for meson production [arXiv:0806.1449]
- Normalization based on the
   MB result [arXiv:0806.1449]
- Focused and aligned to the beam axis if  $\theta \in (30,210)$  mrad
- Meson decay probability (within 50 m) according to the decay law
- Three-body decay of mesons described by MadGraph

- DM upscattering in the detector fiducial volume
- Energy threshold 140 MeV, angular separation  $< 10^{\circ}$ , detection efficiency [arXiv:1407.6060]

#### **Example Fit**

#### Example fit in the double-mediator scenario



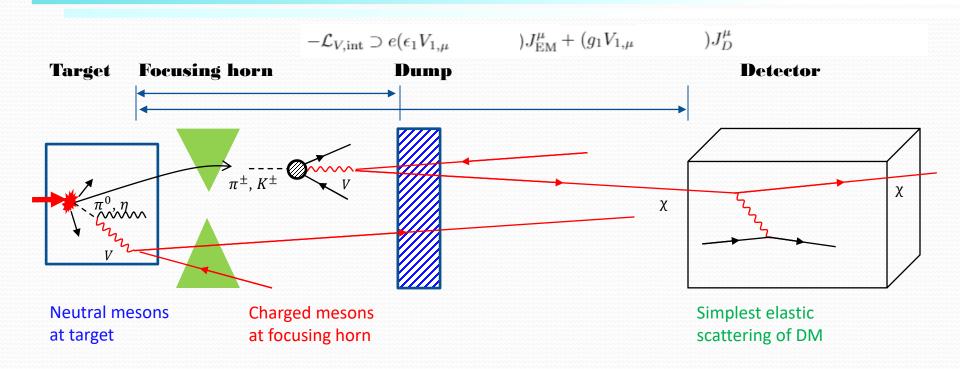
[Dutta, **DK**, Thompson, Thornton, Van de Water, arXiv:2110.11944]

	Vector-portal dark	k matter			
Scenario	$(m_{V_1}, m_{V_2}, m_{\chi}, m_{\chi'})$	$\epsilon_1 \epsilon_2 g_2^{\prime 2}/(4\pi)$	$\chi^2/{ m dof}$		
Single	$(17, -, 8, 40) \; { m MeV}$	$3.6 \times 10^{-9}$	2.5	4	Reasonable fit in the single-mediator scenario
Double	(17,200,8,50)MeV	$1.3 \times 10^{-7}$	2.2		

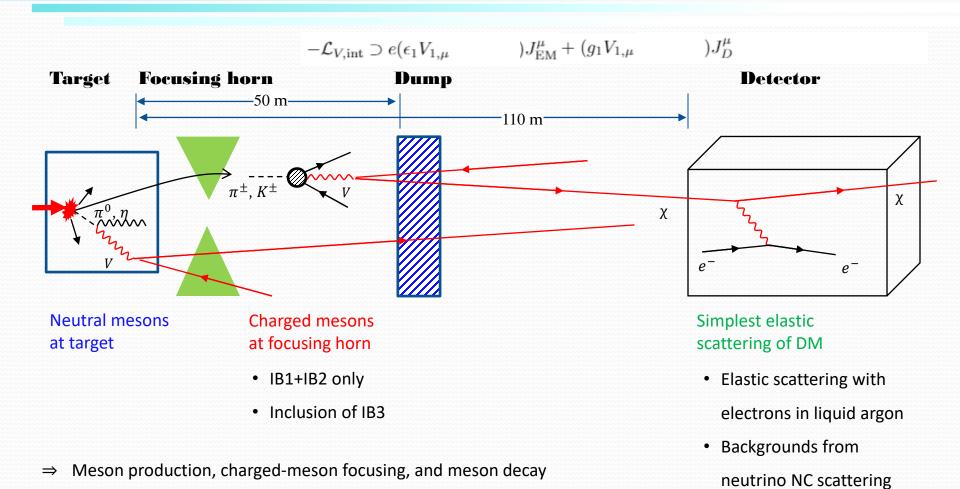
- Best-fit parameter values consistent with the existing constraints including exotic meson decays, dark photon limits, MB off-target mode measurements (see the bonus slides)
- Other dark-sector solutions [e.g., long-lived (pseudo)scalar] available (see the bonus slides)

# Example Physics Case 2: Light Dark Matter Search at SBND

#### Vector-Portal Dark Matter Searches in v Experiments: Main Idea



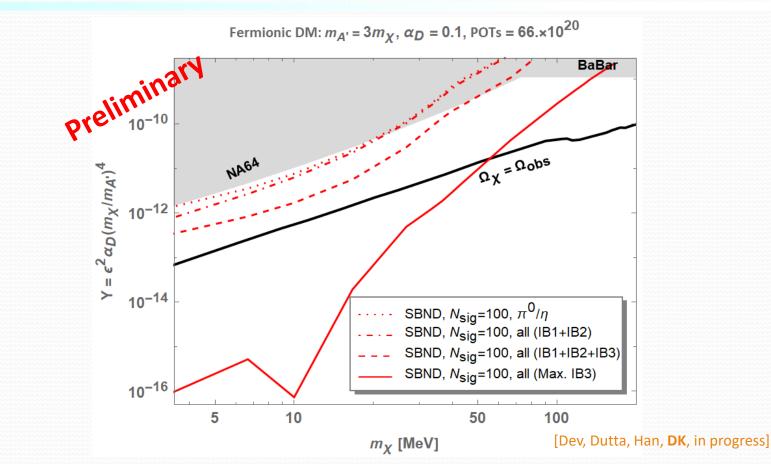
# **Example Data Analysis at SBND**



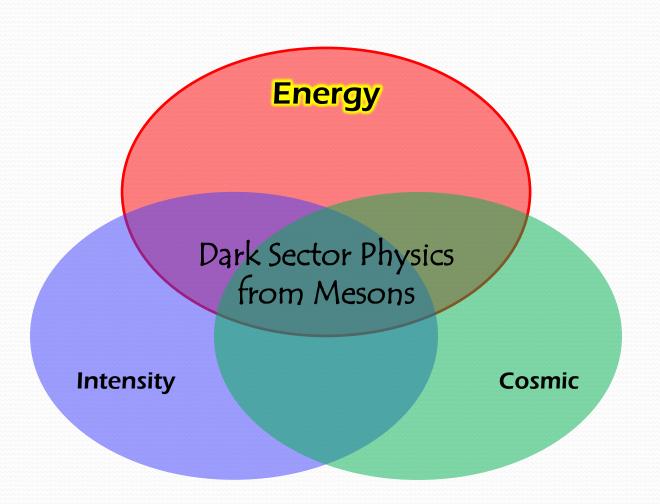
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dynamics simulation schemes similar to those of the MB case

#### **Expected Sensitivity**

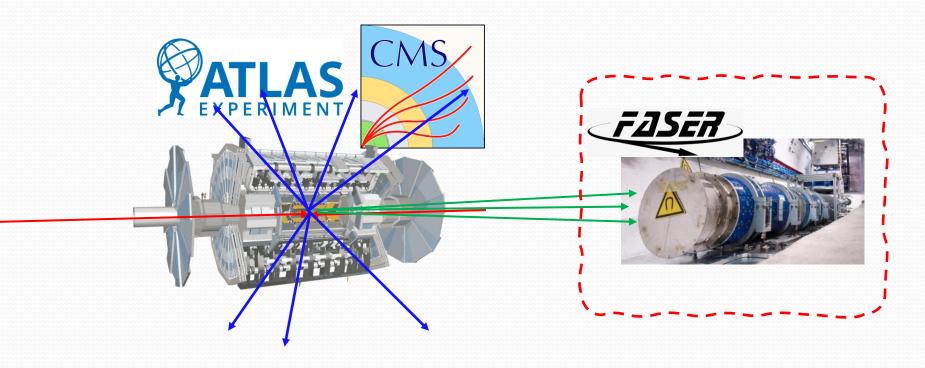


Cf. Maximum IB3 contributions are set by the measurements of exotic charged meson decays:  $\pi \to \ell \nu_\ell A'$  [PIENU Collaboration, arXiv:2101.07381],  $K \to \mu \nu_\mu A'$  [NA62 Collaboration, arXiv:2101.12304].

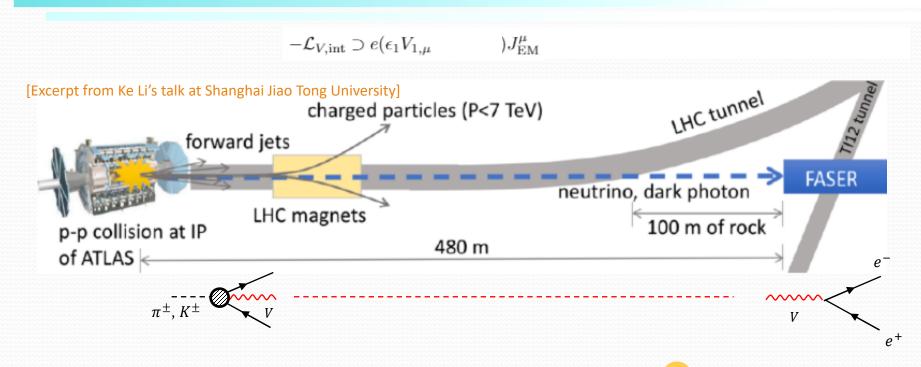


# Example Physics Case 3: Light Mediator Search at LHC

# "Forward" Physics and "Central" Physics

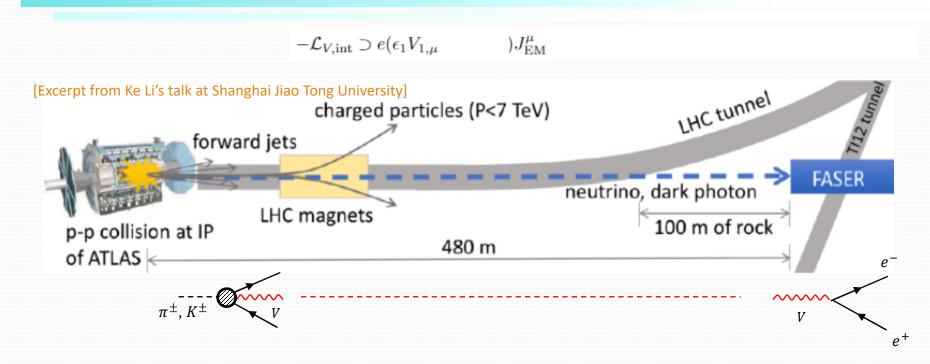


#### Mediator Searches: Main Idea



- Charged mesons are highly boosted, hence effectively focused.
- Distance between the interaction point and the magnet is small, hence a small fraction of charged meson would decay.
- IB3 will enhance the BRs of charged meson exotic decays.

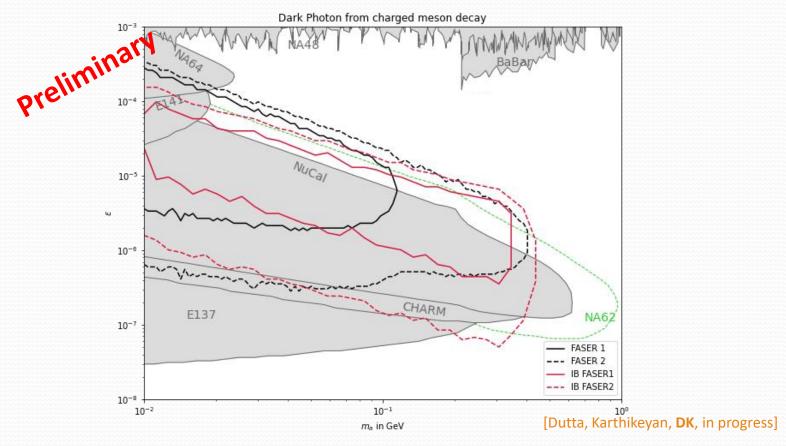
# Data Analysis for the Standard Dark Photon Scenario



- Production of mesons
   described by FORESEE (via
   EPOS-LHC) [Kling, Trojanowski,
   arXiv:2105.07077]
- Charged mesons allowed to decay (conservatively)
   within ~20 m
- IB3 included

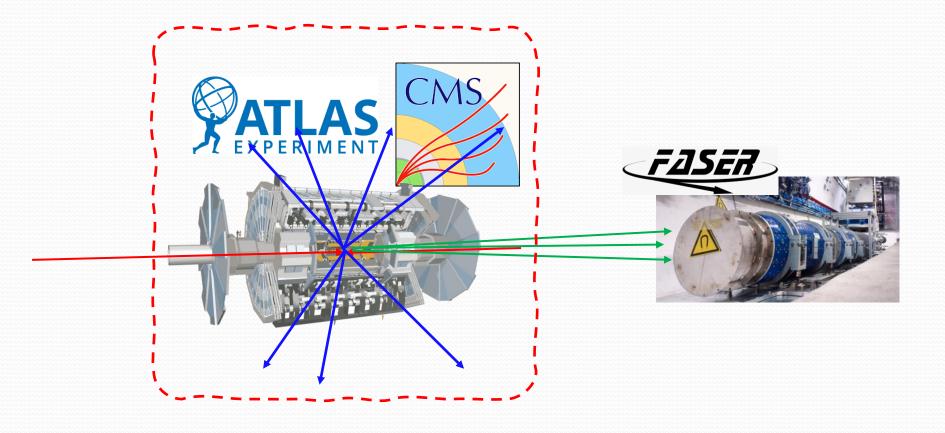
 Produced dark photon should survive before reaching the detector and decay inside the detector.

#### **Expected Sensitivity**

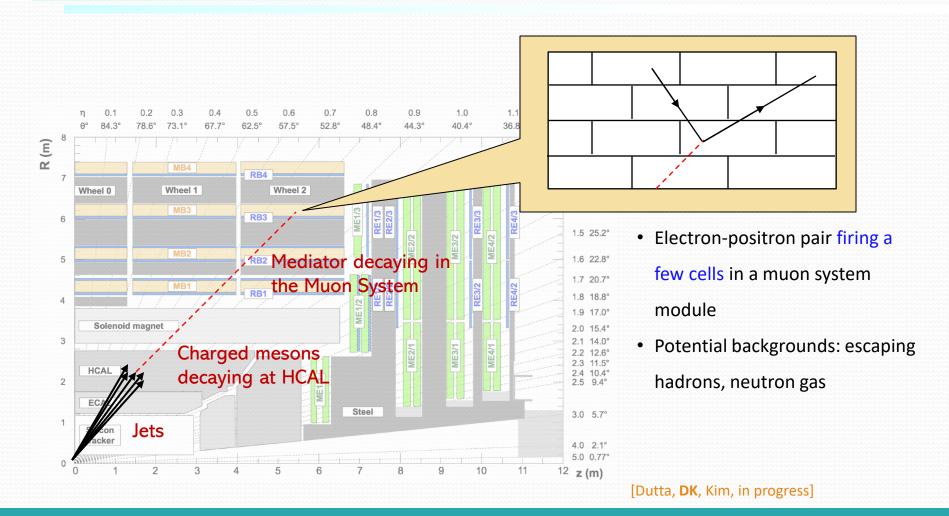


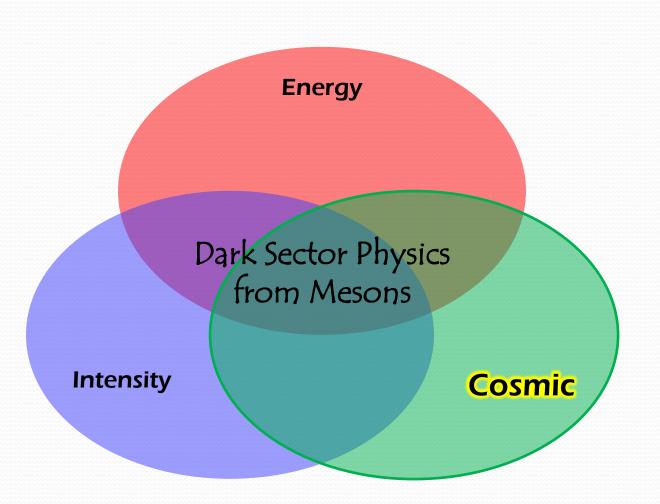
IB3 contributions are maximized with BRs set by the measurements of exotic charged meson decays:  $\pi \to \ell \nu_\ell A'$  [PIENU Collaboration, arXiv:2101.07381],  $K \to \mu \nu_\mu A'$  [NA62 Collaboration, arXiv:2101.12304].

# "Forward" Physics and "Central" Physics

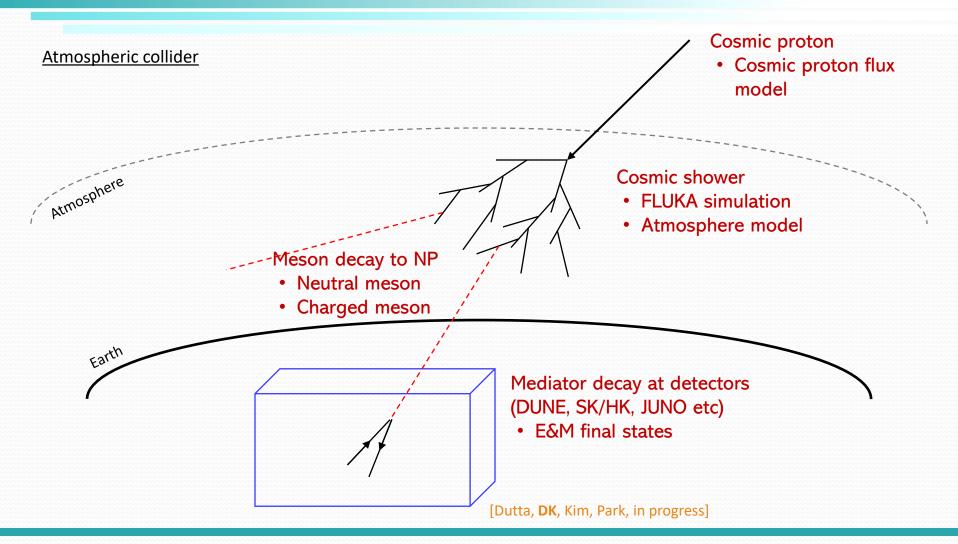


#### Mediator Searches at Muon System: Main Idea

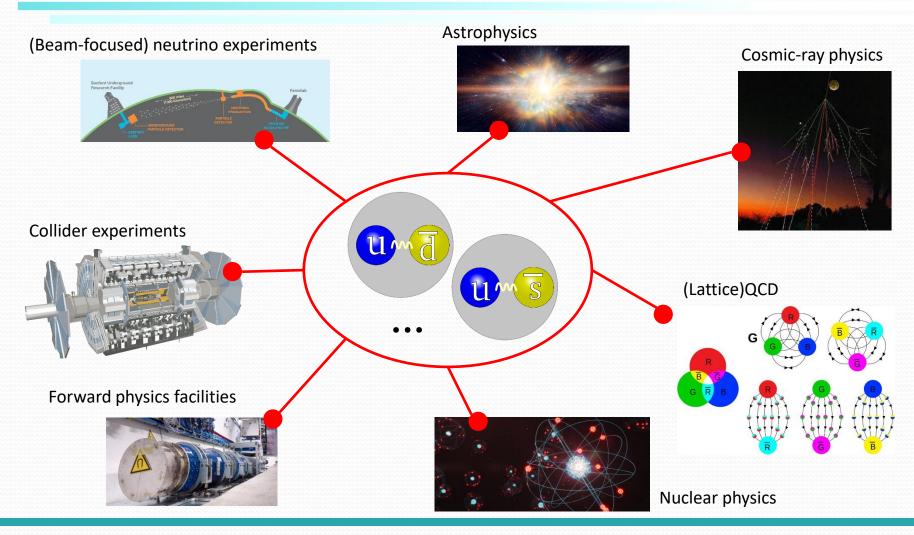




# Dark-Sector Probes Using Cosmic Showers: Main Idea



# Various Physics Opportunities with Charged Mesons



#### **Conclusions**

- Mesons can serve as windows to look into new physics.
- ☐ Charged mesons are overlooked but **efficient sources** of new physics particles (e.g., mediators, dark matter)
  - ✓ BRs of three-body meson decays ≫ BRs of two-body meson decays
  - ✓ Focused flux of new particles (in beam-focused neutrino experiments, forward physics facilities, cosmic-ray experiments)
- □ Dark-sector signals from charged/neutral mesons can be probed in intensity-, energy-, and cosmic-frontier experiments/facilities which provide complementary information.
- Many opportunities are still unexplored.

Thank you!

# **Bonus Slides**

#### **Example UV Models**

- ☐ The mediators (scalar, pseudoscalar, and vector) can emerge from various UV complete models where the scalar/pseudoscalar can couple to the SM particles via Yukawa and the vector mediators via gauge couplings.
- One can use a vector mediator with coupling  $g \sim 10^{-5}$  emerging from pion and then scattering off the nucleus via exchanging a  $\pi^0$  can fit the MiniBooNE excess. The exchanged particle can be a scalar/pseudoscalar as well. Example models include  $U(1)_{T3R}$  [Dutta, Ghosh, Kumar, 1905.02692],  $U(1)_X$  [Correia, Fajfer, 1905.03867].

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#### **Constraints**

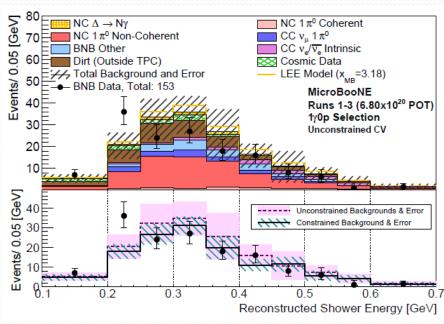
- Limits from exotic decays of charged mesons, e.g., BR( $\pi/K \rightarrow \ell \nu + \nu \nu$  or ee) <  $\sim 10^{-9} 10^{-6}$
- ☐ Limits of visibly-decaying dark photons (e.g., E141, BaBar) and invisibly-decaying dark photons (e.g., NA64)

■ No significant events in the MB off-target mode

We can find respective coupling values that are consistent with these limits.

 $\pi^0$  contribution ~4(8)%, yielding ~2.4  $\pm$  0.5 (4.4  $\pm$  0.9) events in the MB off-target mode for the single(double)-mediator scenario  $\rightarrow$  Consistent with the measurement within the measurement uncertainty.

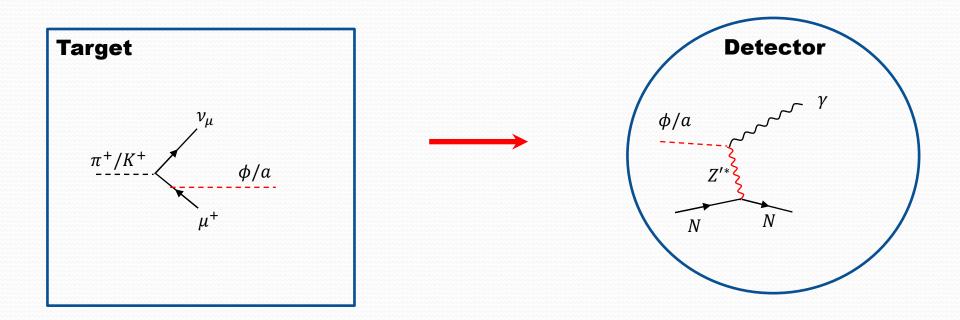
#### MicroBooNE Results and the MB Excess



[MicroBooNE Collaboration, arXiv:2006.16883]

- The (earlier) MicroBooNE result constrains NC  $\Delta \rightarrow N\gamma$  event rates more stringently, supporting that the MB excess requires a new physics interpretation!
- ☐ The later result shows no excess in the channels of a single electron + something.
- ☐ The MicroBooNE data may not be sensitive enough to the MB excess, yet, because of (~3 times) smaller POTs, (~6 times) smaller detector volume, (~3 times) smaller photon efficiency.

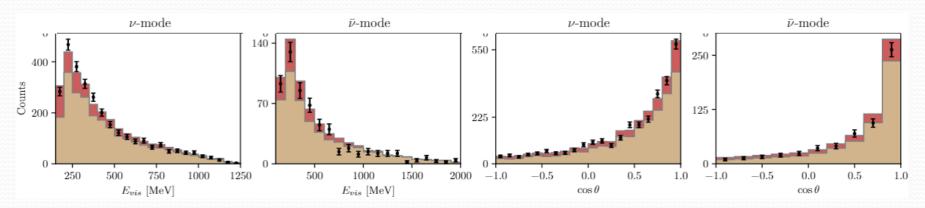
# Example Long-Lived (Pseudo)scalar Models



$$\mathcal{L}_{S(P)} \supset g_{\mu}\phi(ia)\bar{\mu}\mu + g_n Z_{\alpha}'\bar{u}\gamma^{\alpha}u + \frac{\lambda}{4}\phi(a)F_{\mu\nu}'F^{\mu\nu} + \text{h.c.}$$

# **Example Fit**

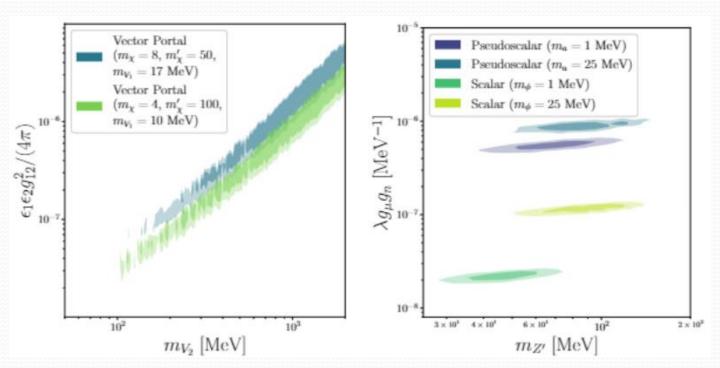
#### Example fit in the scalar scenario



[Dutta, DK, Thompson, Thornton, Van de Water, arXiv:2110.11944]

Long-lived (pseudo)scalar						
Scenario	$(m_{Z'}, m_{\phi/a})$	$(g_{\mu}g_n\lambda) [\text{MeV}^{-1}]$	$\chi^2/{ m dof}$			
Scalar	(49,1) MeV	$2.2 \times 10^{-8}$	1.6			
Pseudoscalar	(85,1) MeV	$5.9 \times 10^{-7}$	1.6			

#### **Credible Regions**



[Dutta, DK, Thompson, Thornton, Van de Water, arXiv:2110.11944]